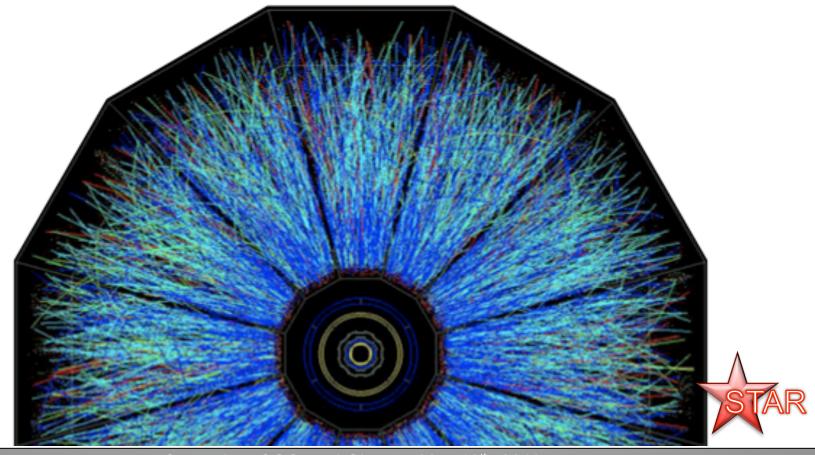
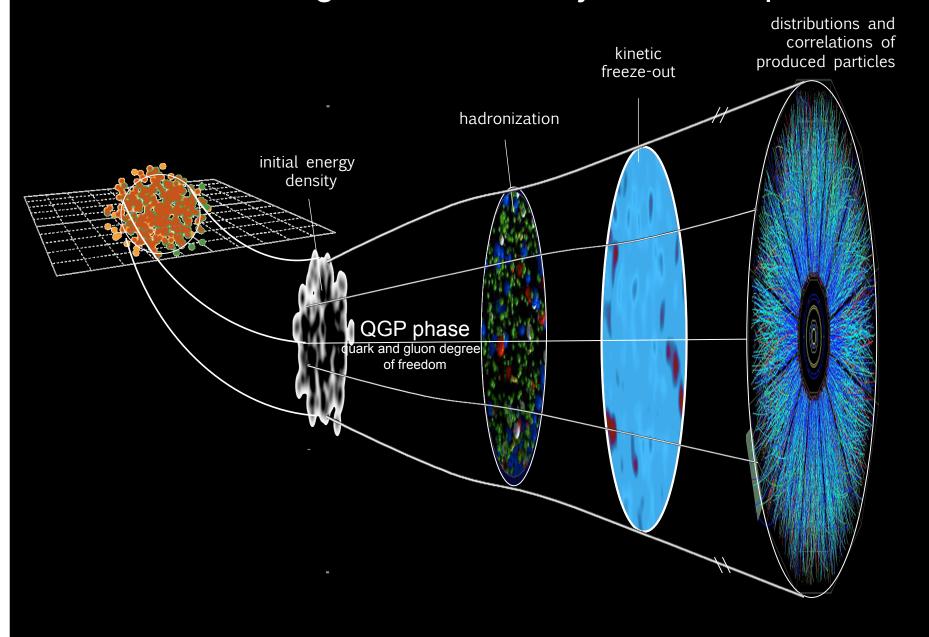
v<sub>2</sub> Cumulant Measurements and Eccentricity Fluctuations for the STAR Collaboration



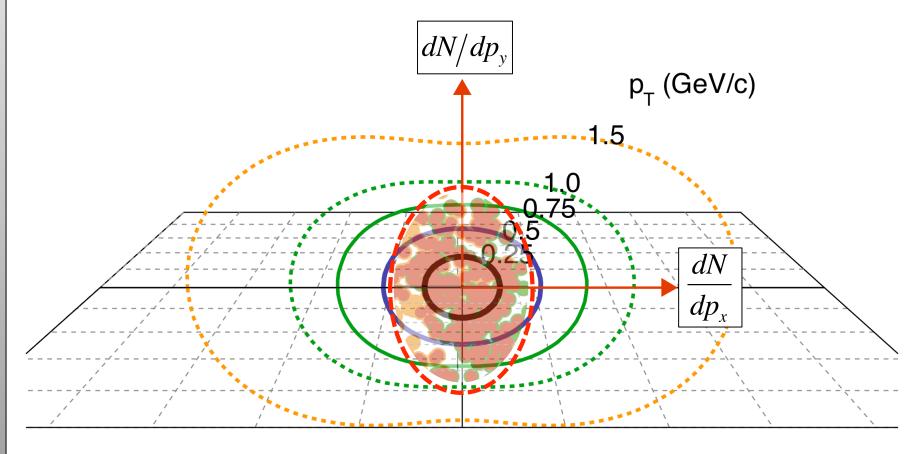


#### Schematic Diagram of the Conjectured Expansion



#### **Azimuthal Distributions**

Collision of two Lorentz contracted Gold nuclei



Are particles emitted at random angles?

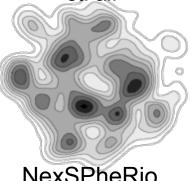
No. They remember the initial geometry.

## **Geometry Fluctuations**

And the initial geometry can be complex.

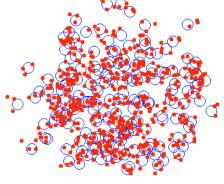
v<sub>2</sub> fluctuations from eccentricity fluctuations will lead to a difference between  $v_2$ {2} and  $v_2$ {4}

Hama, Grasi, Kodama, et. al.



NexSPheRio

Kumar Pruthi, Sorensen



**Additive Quark Model** 

Kowalski, Lappi, Venugopala



IPsat GCG, Glasma

I'll present STAR data on  $v_2\{2\}$  and  $v_2\{4\}$  and compare that data to models for the initial eccentricity

#### v<sub>2</sub> Fluctuations Test Initial Conditions

Data show  $v_2$  depends on eccentricity:  $v_2 = c^* \varepsilon$  where c can depend on dN/dy,  $\sqrt{snn}$ , etc.

$$\sigma_{v2}^{2} \approx c^{2}\sigma_{\epsilon}^{2} + \epsilon^{2}\sigma_{c}^{2} + \text{cross-terms}$$

When we compare to initial eccentricity models we will neglect  $\sigma_c$  so that

$$\sigma_{v2}/v_2 = \sigma_{\epsilon}/\epsilon$$

Vogel, Torrieri, and Bleicher argue that  $\varepsilon^2 \sigma_c^2$  ( $\Delta_{dyn}^2$ ) is proportional to the Knudsen number (nucl-th/0703031) so  $\sigma_c^2=0$  is equivalent to assuming zero viscosity

## Relationship of $\sigma_{v2}$ to $v_2\{2\}$ and $v_2\{4\}$

$$\sigma_{v_n}^2 = \langle v_n v_n \rangle - \langle v_n \rangle \langle v_n \rangle$$

$$v_2^2\{2\} = \langle \cos(2(\varphi_1 - \varphi_2)) \rangle = \langle v_2 \rangle^2 + \sigma_{v_2}^2 + \delta_{non-RP}$$

$$v_2^2 \{4\} = \sqrt{2v_2^2 v_2^2 - v_2^4} \approx \langle v_2 \rangle^2 - \sigma_{v_2}^2$$

$$v_2^2\{2\} - v_2^2\{4\} \equiv \sigma_{tot}^2 \approx \delta_2 + 2\sigma_{v_2}^2$$

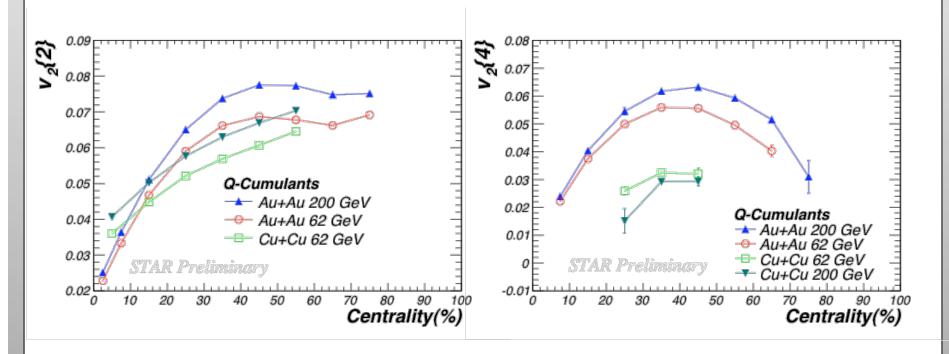
Ollitrault, Poskanzer, Voloshin: Nucl.Phys.A830:279C-282C,2009

Eccentricity fluctuations should show up in the difference between  $v_2\{2\}$  and  $v_2\{4\}$ : but so should non-flow correlations

Note that non-flow is defined relative to either the reaction- or participant-plane

$$\delta_n = v_n^2 \{2\} - \left\langle \cos(n(\varphi - \psi))^2 \right\rangle$$

## STAR Data at 2 Energies and 2 Systems

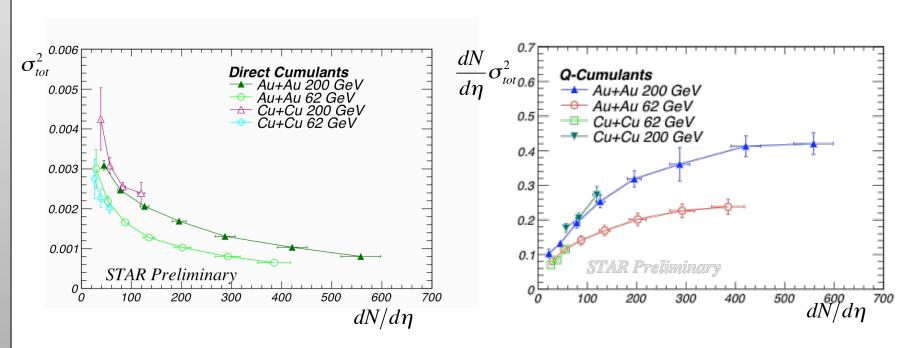


v<sub>2</sub>{2} and v<sub>2</sub>{4} have been measured by STAR for Au+Au and Cu+Cu collisions at 62.4 and 200 GeV

Direct Q-cumulant calculation is used Priv. Com.: Voloshin, Bilandzic, Snellings

We will study 
$$\sigma_{tot}^2 = v_2\{2\}^2 - v_2\{4\}^2$$

# The total width $\sigma_{tot}^2 = v_2\{2\}^2 - v_2\{4\}^2$



Width falls with multiplicity but deviates from the 1/N expected for dilution of correlations with increased combinatorics

Width scales smoothly from Cu+Cu to Au+Au when plotted vs dN/dη

Width scaled by dN/dη increases with centrality (violating a simple linear superposition model for correlations).

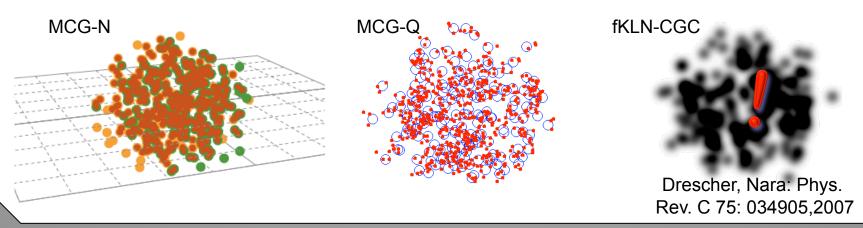
#### **Eccentricity Models**

We will compare this width to the widths predicted from three eccentricity models

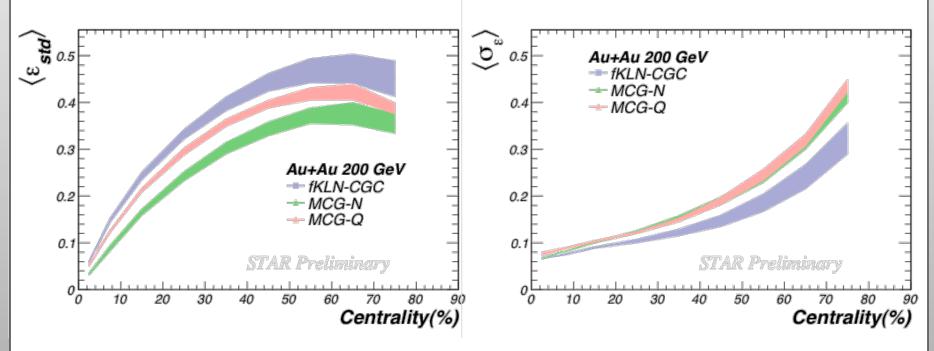
Model 1: Monte Carlo Glauber with nucleons treated as the participants (MCG-N)

Model 2: Monte Carlo Glauber with constituent quarks treated as the participants (MCG-Q)

Model 3: factorized Kharzeev-Levin-Nardi Color Glass Condensate model (fKLN-CGC)



#### Model Results Au+Au



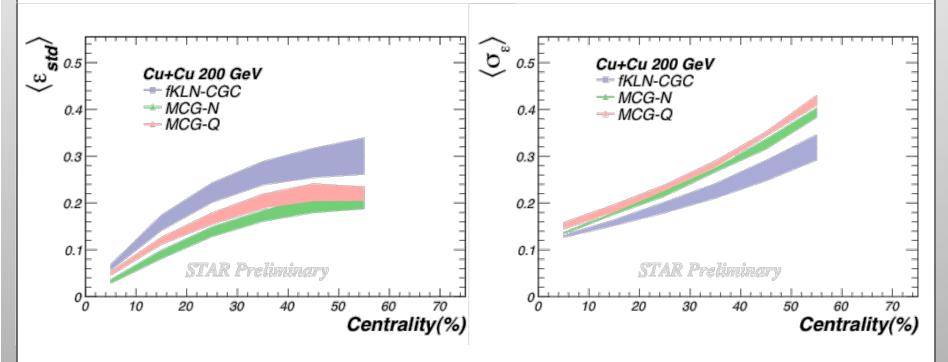
Centrality bins are defined according to the multiplicity from the model: multiplicity modeled using 2-component model ( $x_{hard}$ =0.11)

for eccentricity: fKLN-CGC > MCG-Q > MCG-N

for fluctuations in eccentricity: MCG-Q >~ MCG-N > fKLN-CGC

for  $\sigma_{\epsilon}/\epsilon$ : MCG-N > MCG-Q > fKLN-CGC

#### Model Results Cu+Cu



The same trends hold for Cu+Cu collisions

for eccentricity: fKLN-CGC > MCG-Q > MCG-N

for fluctuations in eccentricity: MCG-Q > MCG-N > fKLN-CGC

for  $\sigma_{\epsilon}/\epsilon$ : MCG-N > MCG-Q > fKLN-CGC

#### **Comparing Data to Models**

 $v_2$ {2} and  $v_2$ {4} provide powerful discrimination between models. In the following slides we'll compare data:

$$\sqrt{\frac{v_2\{2\}^2 - v_2\{4\}^2}{v_2\{2\}^2 + v_2\{4\}^2}} = \frac{\sigma_{v_2}}{v_2} \sqrt{\frac{1 + \delta_2/2\sigma_{v_2}^2}{1 + \delta_2/2v_2^2}} = \frac{\sigma_{v_2}}{v_2} \bigg|_{\text{max}}$$

to model results for:

$$\frac{\sigma_{\varepsilon}}{\varepsilon} = \sqrt{\frac{\varepsilon^2 \{2\} - \varepsilon^2 \{4\}}{2\varepsilon^2 \{4\}}}$$

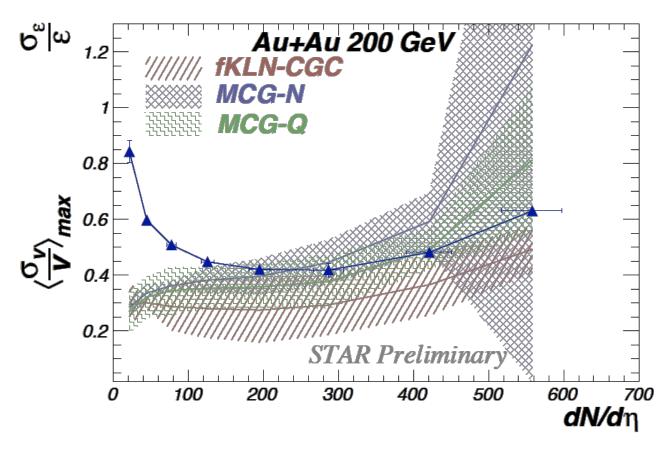
Bhalerao and Ollitrault: Phys.Lett.B641:260-264,2006

In case that non-flow dominates the width

$$\sqrt{\frac{v_2\{2\}^2 - v_2\{4\}^2}{v_2\{2\}^2 + v_2\{4\}^2}} = \sqrt{\frac{1}{1 + 2v_2^2/\delta_2}}$$

Ratio is 1 if non-flow dominates and  $v_2$ =0 or  $\sqrt{(4/\pi-1)}$  if  $\epsilon$  fluctuations dominate

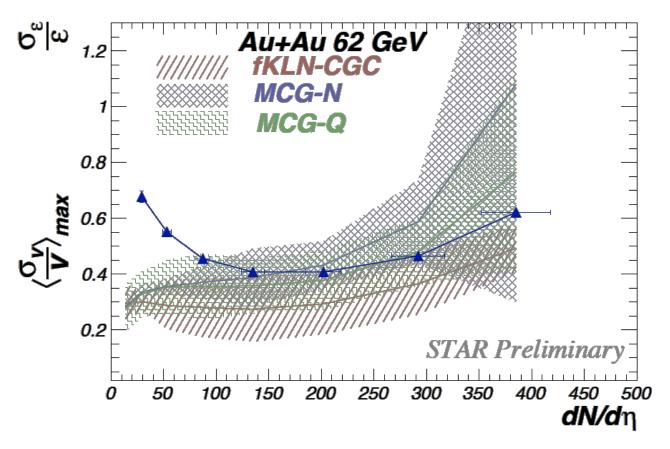
## Comparison of models to $\sigma_{tot}$



For central 200 GeV Au+Au collisions, the width expected from MCG-N eccentricity fluctuations nearly exceeds the total width of data

MCG-Q and fKLN-CGC remain smaller and consistent with  $\delta_2$ >0

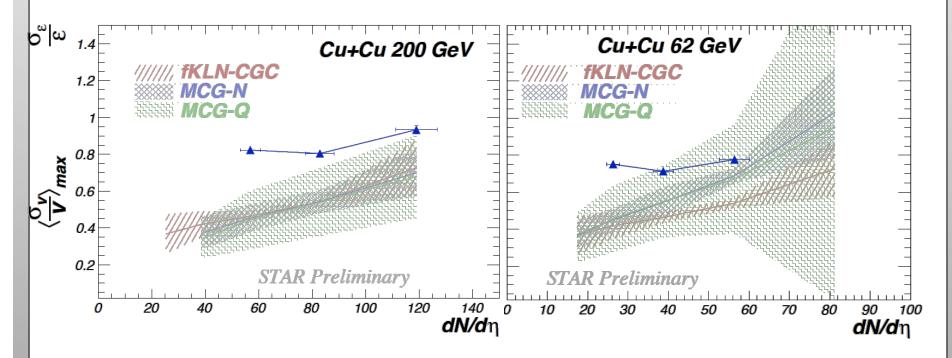
## Comparison of models to $\sigma_{tot}$



For central 62.4 GeV Au+Au collisions, the width expected from MCG-N and MCG-Q ε fluctuations nearly exceeds the total width of data

Only fKLN-CGC remains smaller and consistent with  $\delta_2$ >0

## Comparison of models to $\sigma_{tot}$



For Cu+Cu collisions at both energies, data is wider than all three models

Data are limited by difficulty in determining v<sub>2</sub>{4} at small multiplicities

#### What About Non-flow?

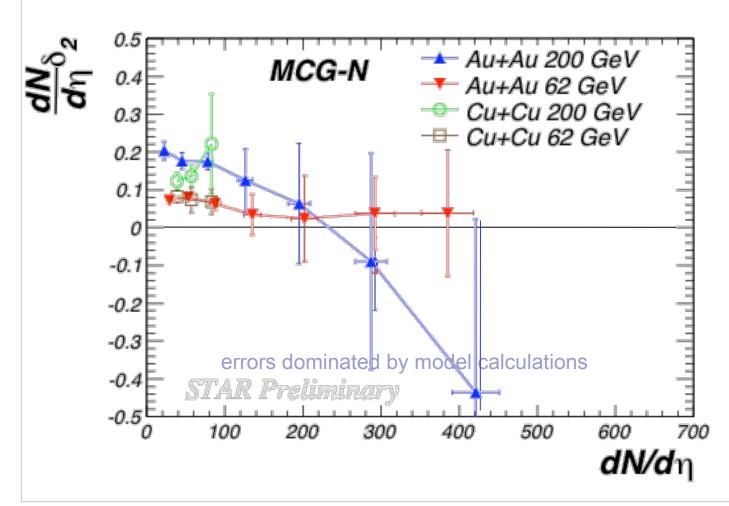
The previous comparisons can be extended by calculating the width that remains after subtracting off the eccentricity fluctuations implied by each model, *e.g.* 

$$\delta_2^{MCG-N} = \sigma_{tot}^2 - 2\left(v_2 \frac{\sigma_{\varepsilon}^{MCG-N}}{\varepsilon^{MCG-N}}\right)^2$$

Let's see what that looks like for each model.

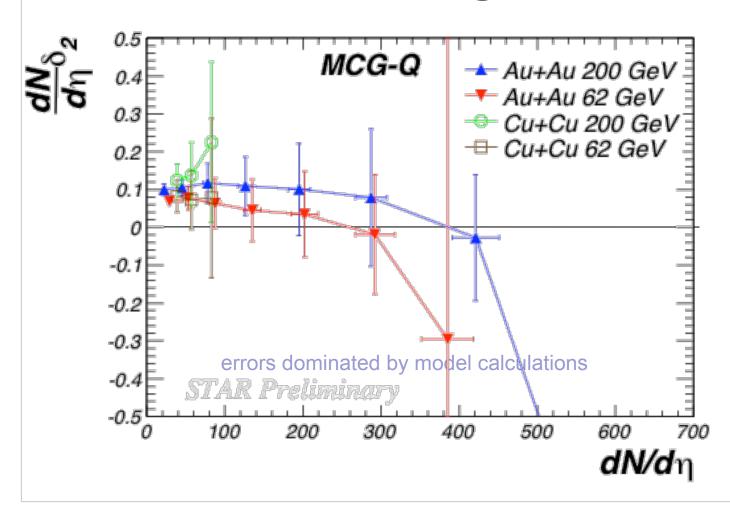
We'll also scale  $\delta_2$  by dN/d $\eta$  to account for dilution of correlations with increased multiplicity

#### The remaining width



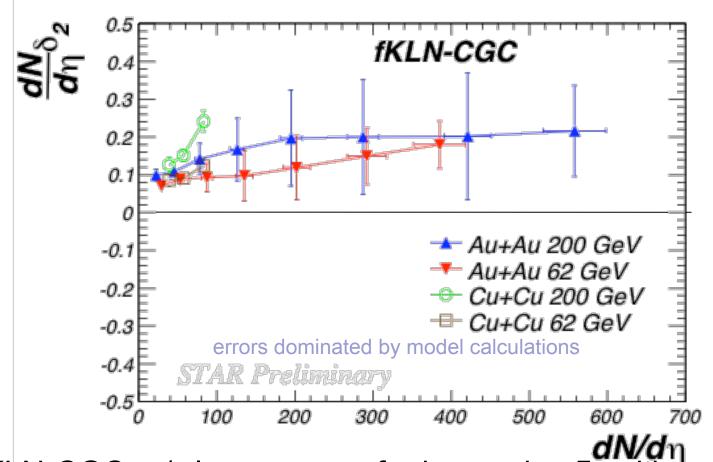
For the MCG-N eccentricity fluctuations to be correct, nonflow would need to be nearly zero or negative in central Au +Au collisions

#### The remaining width



The MCG-Q eccentricity fluctuations do not require negative non-flow in central 200 GeV Au+Au but still do for 62.4 GeV

#### The remaining width



fKLN-CGC  $\sigma_{\epsilon}/\epsilon$  leaves room for increasing  $\delta_2$  with centrality:

 $\sigma_{\epsilon}$  and  $\epsilon$  calculations can be supplemented with predictions for  $\delta_2$  to check for consistency

## v<sub>2</sub>/ε scaling

Now we can address  $v_2/\epsilon$  scaling in a consistent way:

 $v_2$  is determined assuming  $v_2$  fluctuations as predicted by the eccentricity fluctuations of each model

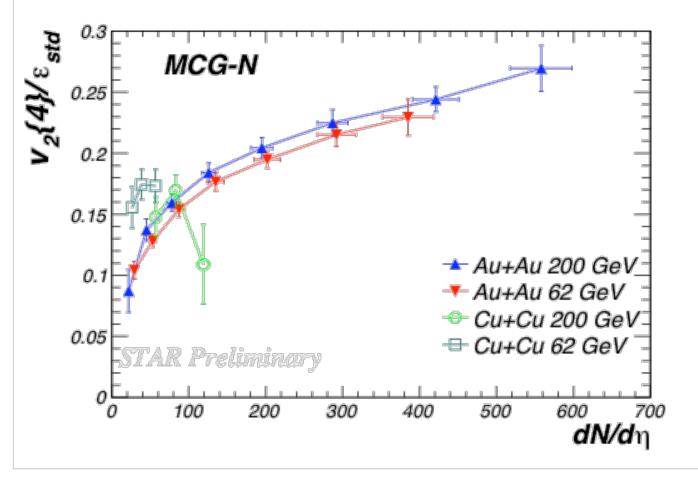
the eccentricity is calculated from the same model

It turns out that in the case that eccentricity fluctuations dominate  $v_2$  fluctuations, this reduces to:

$$\frac{\langle v_2 \rangle}{\langle \varepsilon \rangle} = \frac{v_2 \{4\}}{\varepsilon_{std}}$$

where  $\varepsilon_{std}$  is the eccentricity calculated relative to the reaction plane not the participant plane.

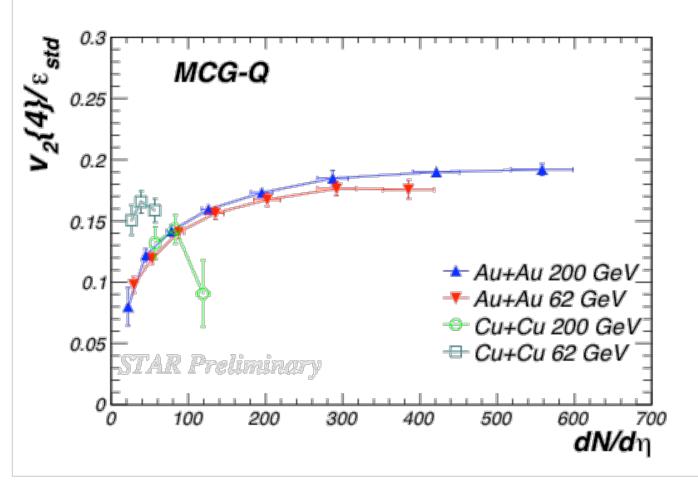




For the MCG-N model,  $v_2/\epsilon$  rises continuously

No indication of a saturation at a hydro-limit

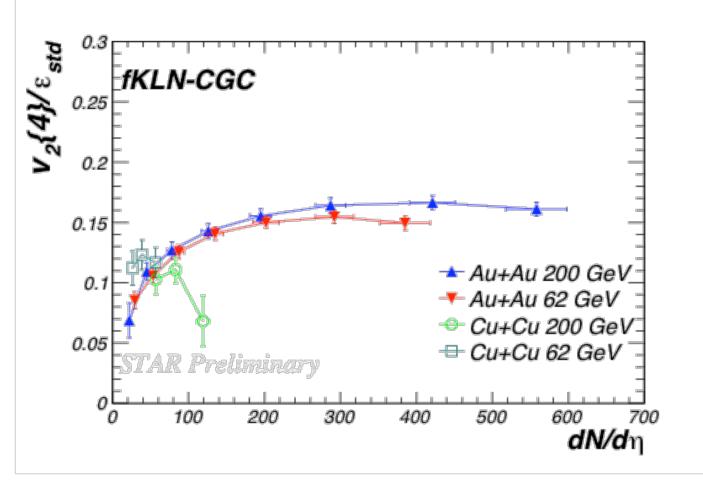
## MCG-Q: $v_2/\epsilon$ Scaling



For the MCG-Q model,  $v_2/\epsilon$  rises then starts to level-off

Only a small increase in  $v_2/\epsilon$  for events with  $dN_{ch}/d\eta > 300$ 

## fKLN-CGC: v<sub>2</sub>/ε Scaling



For the fKLN-CGC model,  $v_2/\epsilon$  rises then saturates

For  $dN_{ch}/d\eta > 250$ ,  $v_2$  scales with  $\varepsilon$ 

#### **Conclusions**

The 2- and 4-particle v<sub>2</sub> cumulants have been measured for Au+Au and Cu+Cu collisions at 62.4 and 200 GeV

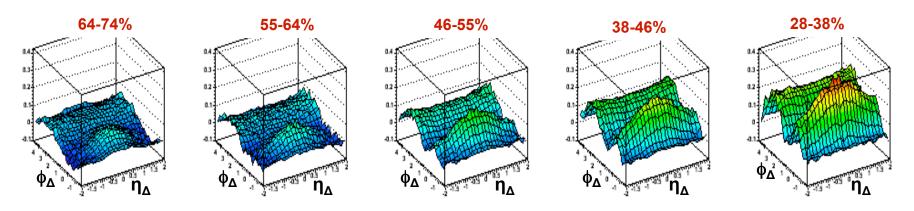
We used the difference  $v_2^2\{2\}-v_2^2\{4\}$  to test models of the initial eccentricity (the difference is a measurement not an error)

MCG models predict larger eccentricity fluctuations in central Au+Au collisions leaving little room for non-flow effects while the fKLN-CGC model is well within the range allowed by  $\sigma_{tot}$ 

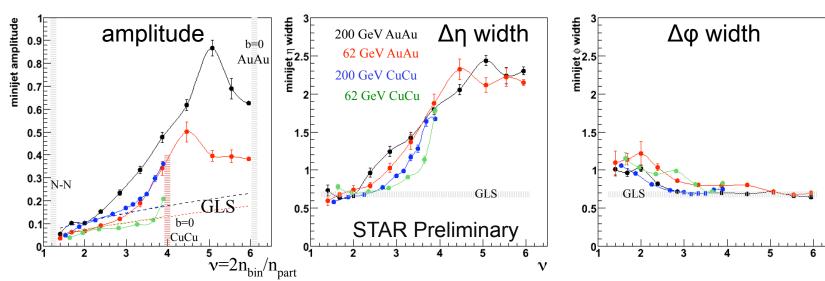
Above dN/dη~200, v<sub>2</sub> scales with fKLN-CGC eccentricity but not MCG-N eccentricity

For discussion of 2-particle correlations relevant to non-flow see Lanny Ray's talk later this week

#### 2-Particle Correlations



Correlations Between All Pairs: HBT, and photon conversion pairs subtracted



## Ridge and Cone Phenomenology

Chemical composition of the ridge & cone

 $\triangleright$  Baryon-to-Meson ratios like the bulk (p/ $\pi$  and K<sub>S</sub>/ $\Lambda$ )

#### Correlation amplitude

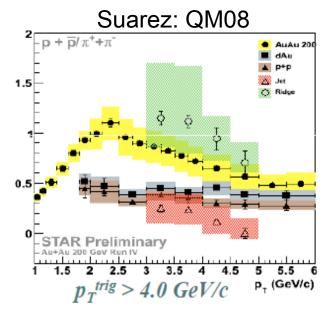
- ➤ Correlations increase faster than N<sub>bin</sub> or N<sub>part</sub>; closer to M(M-1) instead
- ➤ Near and Away-side amplitudes have same centrality dependence

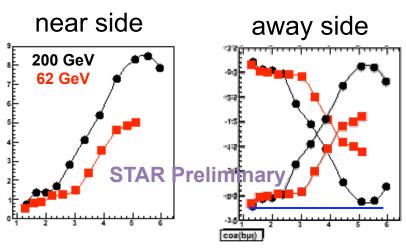
#### Longitudinal and Azimuthal Width

> both different from fragmentation

#### p<sub>T</sub> spectra of the ridge and cone

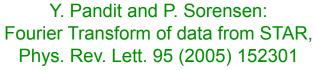
➤ Both are soft; like the bulk not like jet fragments

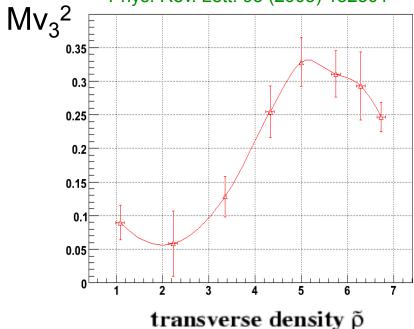




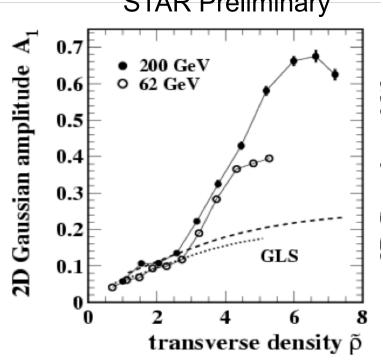
Lanny Ray: CATHIE RIKEN workshop

#### What's So Odd About the Ridge and Cone?





#### low p<sub>T</sub> ridge yield STAR Preliminary



Large possible  $\langle v_3^2 \rangle$  component in intermediate  $p_T$  data

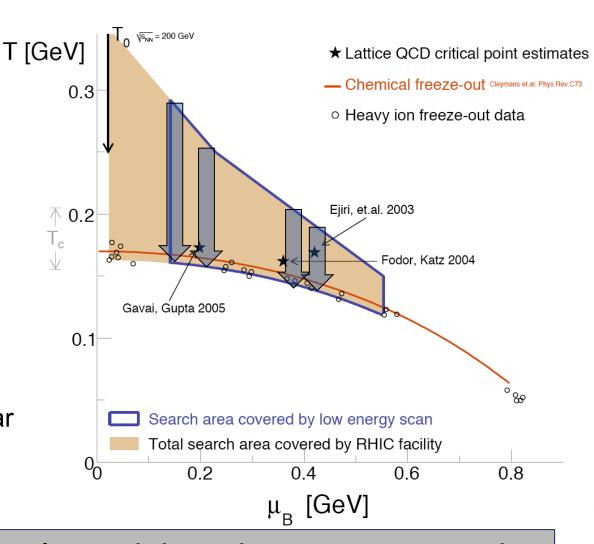
Centrality dependence is similar to the low p<sub>T</sub> ridge

#### Search for a critical point at RHIC



In 1911, Rutherford discovered the nucleus, making him the first nuclear physicist

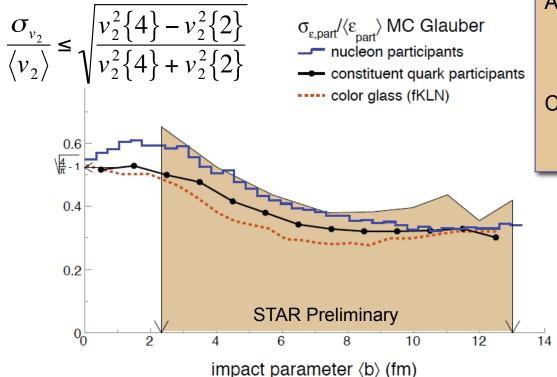
100 years later, RHIC will scan for new landmarks on the nuclear matter phase diagram



The experimental search is underway as we speak

## **Comparison to Models**

# Upper limit challenges models: MC Glauber already exhausts entire width with participant fluctuations



#### Additive Quark MC:

treats confined constituent quarks as the participants decreases eccen. fluctuations

#### Color Glass MC:

includes effects of saturation increases the mean eccentricity

comparison to hydro (NexSPheRio): *Hama* et.al. arXiv:0711.4544

eccentricity fluctuations from CGC: *Drescher, Nara. Phys.Rev.C76:041903,2007* 

extraction of Knudsen number: *Vogel, Torrieri, Bleicher, nucl-th/0703031* 

fluctuating initial conditions: *Broniowski*, *Bozek*, *Rybczynski*. *Phys.Rev.C76:054905*,2007

first disagreement with  $\varepsilon_{standard}$  and use of quark MC: *Miller, Snellings. nucl-ex/0312008*